OPTICAL AND MORPHOLOGICAL INVESTIGATIONS OF CHALCOGENIDE GE-SB-TE THIN FILMS

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Abstract. Chalcogenide thin films based on Ge-Sb-Te (GST) are synthesized from elements with 5N purity using the conventional melt-quenching method. The mixture was used as material for deposition on quartz substrate by the chemical vacuum deposition method. Optical and dielectric properties were determined by Spectro-ellipsometry measurements. The optical dispersion parameters are calculated using the single-oscillator model of Wemple-DiDomenico. The dissipation factor, the optical conductivity, the complex electric modulus, and the complex impedance of these GST layers were calculated. The AFM-SEM imagines revealed compact, continuous, and smooth films with good adhesion and uniform granules distribution.

Keywords: chalcogenide layer, thin film, dielectric properties, optical conductivity.

1. INTRODUCTION

Chalcogenide thin films based on Ge-Sb-Te (GST) are intensely studied due to optical, dielectric, and mechanical structural properties [1-3]. These materials are identified by a high enough contrast in electrical resistivity or optical phase-change memories parameters [4]. The changes from the amorphous state with high electrical resistivity to crystalline structure with low electrical resistivity are quantified by the Joule effect using an electric current pulse [5]. GST is an innovative material with high structural quality for phase change, offered by epitaxial films. Progress is made in the growth of crystalline phase change materials by physical methods, such as molecular beam epitaxy, sputtering, and pulsed laser deposition the difference and similarities between these physical deposition methods and the crystal structures of $Ge_2Sb_2Te_5$.

Over time a few methods were developed for chalcogenide layers, such as spin coating [6], magnetron sputtering [7, 8], thermal evaporation [9], atomic layer deposition [10], and metal-organic vapor phase epitaxy [11]. The Pulsed Laser Deposition (PLD) technique is suitable for the thin film growth of complex materials with good adhesion to the substrate and high homogeneity [12–18]. Experimentally technique discussed in above mentioned articles is important for the thin film growth because it show the reduction of cross-talk by greater than 30 times and the bending loss by greater than 3 times in densely integrated, ultra-compact photonic circuit blocks. By using a prototype, the dielectric metamaterial-waveguide property achieves a low propagation loss of approximately 3.7 ± 1.0 dB/cm, comparable to those of silicon strip waveguides [12–14].

Growth of crystalline phase change materials by physical deposition methods is discussed in [19] and the phase-change materials for rewriteable data storage in [20]. Amorphous films such GeTe-Sb₂Te₃ and GeSb₂Te₄ or Ge₂Sb₂Te₅ present a large optical change and high-speed data rewriting which are important characteristics of optical memory material [21]. The growth depends on the electrical conductivity, surface morphology, crystallinity, and texture of the film. GST deposited on glasses substrate TiN and TiO₂ are highergrowth substrates, SiO2, Si₃N₄, and ZrO₂ are lower-growth substrate and HfO₂ is a non-growth substrate with no film growth [28, 29]. It is usually expected that the substrate dependency of the growth rate disappears after surface reconstruction is masked by a thicker film [22].

The optical and electrical properties of Ge-Sb-Te thin layers are compared between thermal evaporation (TE) and laser ablation. The ratio Sb/Te from the stoichiometric value in the TE samples [5] revealed a slight variation in chemical composition analysis.

The optical constants (refractive index and extinction coefficient) of thin films deposited by the PLD method presents higher values than those observed for the TE methods. High refractive index materials are commonly integrated with materials of contrasting refractivity to form transparent components with anti-reflective properties.

In this paper, the optical properties of the $Ge_xSb_{20-x}Te_{80}$ (GST) synthesized from elements with 5N purity (Ge, Sb, Te) are analyzed by using the Spectro-ellipsometry experimental data. The novelty of the paper consists in new observations on the optical properties and the photo-induced modifications in chalcogenide compositions.

Finally, the optical and morphological properties are discussed in connection with application of amorphous chalcogenides in plasmonic resonance structure configuration, binary chalcogenide photovoltaic cells and IR-coatings.

2. EXPERIMENTAL PROCEDURE

Three chalcogenide systems glassy $Ge_xSb_{20-x}Te_{80}$ (GST) alloys with composition x = 15, 17, and 19 at %, are synthesized from elements with 5N purity (Ge, Sb, Te) using the conventional melt-quenching method [23-32]. The mixture of the elements with proper weight percent was put in a quartz ampoule sealed down to 10^{-3} Pa pressure. The ampoules were heated to a rotary oven to 950°C, and to obtain homogeneous melting, the materials mixture was kept at this temperature for 24 hours by rotating the oven continuously. After finishing the process, the ampoules were removed, and the melts were cooled in ice water. The material mixture was used as material for deposition on quartz substrate by the thermal and chemical vacuum deposition method. Plasma cleaning discharge took place at 10^{-1} Pa, for 12 minutes in the atmosphere of N₂, and during this time, the dome is rotating (about 10 rot/min). This process is important because it removed the impurities on the walls from the vacuum deposition installation and activated the bulk before deposition. When the pressure has stabilized at 4×10^{-3} Pa, the heating resistance is coupled, ensuring a temperature of 300°C in the evaporation chamber while the dome is rotating with a 10rpm rotation speed. The material mixture was evaporated to 270 mA electric current. The optical monitoring control was done by TFCalc 3.5 software for approximately 3 min with a deposition rate of 6 nms⁻¹ and showed 570 nm thickness. After evaporation, the samples were measured by the Spectro-Ellipsometry method using VASE Ellipsometer-Woollam. The surface quality of the films was characterized by scanning electron microscopy (SEM) and atomic force microscopy (AFM) using technology from ZEISS. The 2D AFM images for chalcogenide systems glassy Ge_xSb_{20-x}Te_{s0} alloys with composition x = 15, 17, and 19 at % were taken from a 5 μ m×5 μ m scanned surface area.

3. RESULTS AND DISCUSSION

Ellipsometry measurements determined optical and dielectric properties. The experimental dates were recorded in the spectral range UV-VIS-NIR, in the spectral range 250–1350 nm, at 50⁰ incidence angles. Fig. 1 shows the experimental dates ψ and Δ for chalcogenide systems glassy Ge_xSb_{20-x}Te₈₀ alloys with composition x = 15, 17, and 19 at %, with a mean square error (MSE) of less than 5 were obtained in the spectral range 250–1100 nm. The model presented in Fig. 2 contains two layers on the quartz substrate: roughness layer /chalcogen layer / substrate.



Fig. 2 – The model layers on the quartz substrate.

The Gauss and Lorentz oscillators simulated the chalcogenide layer by the General Oscillator method [33, 34]. The roughness layer, which is considered a mixture of 50% material (film) and 50% voids (air), was modeled with Bruggeman's effective medium approximation (B-EMA). The optical constants and thickness

249

are determined. The refractive index, n and extinction coefficient, k, according to photon energy, hv are presented in Fig. 3. The thickness for all alloys' chalcogenide composition is around 600 nm.

The absorption coefficient is calculated using the relation



Fig. 1 – The experimental dates ψ and Δ for chalcogenide systems glassy Ge_xSb_{20-x} eso alloys with composition x = 15, 17, and 19 at %: a) ψ and Δ for Ge₁₅Sb₅Te₈₀; b) ψ and Δ for Ge₁₇Sb₃Te₈₀; c) ψ and Δ for Ge₁₉Sb₅Te₈₀.

Figure 4 shows the absorption coefficient versus photon energy, hv. According to Figs. 3 and 4, the constant optical decrease when germanium / antimony content increase / decrease. The relation between the absorption coefficient and the photon energy hv is given by the relation developed by [35, 36]

$$\alpha = \frac{1}{h\nu} \Big[\tilde{A}(h\nu) - E_g \Big]^p, \qquad (2)$$

where \tilde{A} is constant depends on the transition probability, E_g is the width of the optical band gap, and p is an index that characterizes the optical absorption processes in the investigated films. The experimental analysis showed a proportionality between the absorption coefficient and the frequency of the photon energy in the form $(hv-E_g)^p$. The refractive index, the band gap and the oscillator parameters of amorphous GaSe thin films is discussed in [37] and the influence of temperature on charge transports and optical parameters for the Ge₁₅Sb₅Se₈₀ and Ge₁₅Sb₅Te₈₀ amorphous thin films in [38]. The behavior of the electronic dielectric constant in covalent and ionic materials is analyzed in [39]. The experimental results are based on the fundamental ε_2 spectrum, and it is found that for a certain choice of a model ε_2 spectrum, the bandgap parameter E_a in the high-frequency sum rule introduced by Hopfield, provides the connection between the single-oscillator parameters (E_0, E_d) and the static-dielectric-constant parameters ($E_g, \hbar \omega p$), i.e., $(\hbar \omega_p)^2 = E_a E_d$ and $E_g = (E_a E_0)^{1/2}$. Finally, it is suggested that the observed dependence of E_d on the coordination number and valency implies that an understanding of refractive-index behavior may lie in a localized molecular theory of optical transitions.

The exponent p defines the type of the optical transition in the chalcogenide films and should be one of these values: 2, 1/2, 3, and 3/2. Theoretically, p equals to 2, 1/2, 3, or 3/2 for the indirect allowed, direct allowed, indirectly forbidden, and directly forbidden transitions, respectively [40].



Fig. 3 – The optical constants, refractive index, n, and extinction coefficient, k, according to hv.

The usual method [31] for determining the optical transition includes plots of $(\alpha h \upsilon)^{1/p}$ versus the incident photon energy $(h\upsilon)$. A set of plots is given by the exponent $p: (\alpha h \upsilon)^{1/2} - h\upsilon; (\alpha h \upsilon)^{2/3} - h\upsilon; (\alpha h \upsilon)^{1/3} - h\upsilon;$ and $(\alpha h \upsilon)^{2/3} - h\upsilon$. One of these plots determines the type of the optical transition and satisfies the widest linearity

of data in the high absorption region, the exponent p indicates that the dominant transition is a direct allowed one. Consequently, $(\alpha h \upsilon)^2$ was plotted against h υ in Fig. 5.



Fig. 4 – Absorption coefficient versus photon energy, hv.



Fig. 5 – The plotted $(\alpha h\nu)^2$ versus photon energy $h\nu$ for: a) Ge₁₅Sb₅Te₈₀; b) Ge₁₇Sb₃Te₈₀; c) Ge₁₉Sb₁Te₈₀.

Influence of temperature on charge transports and optical parameters for the $Ge_{15}Sb_5Se_{80}$ and $Ge_{15}Sb_5Te_{80}$ amorphous thin films is studied in [32].

 $Ge_xSb_{20-x}Te_{80}$ (GST) alloys with composition x = 15, 17, and 19 at %, were measured by VASE ellipsometer- Woollam. Using the UV-Vis-NIR spectra, the Urbach energy was detected in Fig. 6.



Fig. 6 – The variation of Urbach energy deduced by $ln(\alpha)$ of: a) $Ge_{15}Sb_5Te_{80}$; b) $Ge_{17}Sb_3Te_{80}$; c) $Ge_{19}Sb_1Te_{80}$.

The N_c value is the average coordination number of covalent bonds per atom and was calculated from the relation

$$N_c = [4(x) + 3(40 - x) + 2(60)]/100, \qquad (3)$$

where x is the Ge content in at %, and the numbers 4, 3, and 2 are the valences of the elements Ge, Sb, and Te, respectively, the values of average coordination number N_c are present in Table 1. The variations on optical energy (band gap and Urbach energy) versus Germanium content are in Fig. 7. We can see the values of optical energy decrease when Germanium content increases.



The average coordination number of covalent	
Chalcogenide Layer	Coordination number (N _c)
Ge15Sb5Te80	2.55
Ge17Sb3Te80	2.57

2.59

Table 1

Fig. 7 – The variations on optical energy (Optical band gap E_g and Urbach energy E_u) versus Ge content (%).

Dielectric function. The dielectric function is described by $\tilde{\epsilon} = \epsilon_1 + i\epsilon_2$, where ϵ_1 is the real part that represents the normal dielectric constant, and ϵ_2 is the imaginary part that describes free-carrier absorption of the dielectric constant. The real part ϵ_1 , and imaginary ϵ_2 [25] are described

$$\varepsilon_1 = n^2 - k^2 = \varepsilon_{\infty} - \left(\frac{e^2 N}{4\pi c^2 \varepsilon_0 m^*}\right) \lambda^2 = \varepsilon_{\infty} + 4\pi \chi_p, \qquad (4)$$

Ge₁₉Sb₁Te₈₀

$$\varepsilon_2 = 2nk = \left(\frac{\varepsilon_{\infty}\omega_p^2}{8\pi^2 c^2 \tau}\right)\lambda^3,\tag{5}$$

where ω are the photon angular frequency, τ optical relaxation time, respectively, ε_{∞} is the high-frequency dielectric constant, λ is the wavelength, the free charge carrier concentration, ε_0 , is the permittivity of free space (8.854×10⁻¹² F/m), m^* is the effective mass of the charge carrier, and *c* is the velocity of light, and χ_p is the electric susceptibility. The complex dielectric constant $\tilde{\varepsilon}$ completely describes the reflection, propagation, and light loss in the specimen structure, and therefore the material electronic structure can be clearly described. Fig. 8 represents the real, and the imaginary parts of the dielectric constants studied in photon energy (*h* υ) range Ge_xSb_{20-x}Te₈₀ (GST) alloys with composition *x* = 15, 17, and 19 at %. We can see the values of the real and the imaginary parts of the dielectric constants decrease when Germanium content increase and Antimony content decrease.



Fig. 8 – The real and the imaginary parts of the dielectric constants studied in photon energy $(h\nu)$ range.

According to the oscillating simple model [33, 34], the real part of the permittivity is expressed by:

$$\varepsilon_1(\omega) = 1 + \omega_{p'}^2 \sum_i \frac{f_i}{\omega_i^2 - \omega^2}, \qquad (6)$$

where $\omega_{p'}$ is the angular frequency of plasma, and f_i is the electric force of dipole oscillator for the transitions to the pulsation ω_i .

The expression of the permittivity becomes

$$\varepsilon_{1}(\omega) - 1 = n^{2}(h\nu) - 1 = \frac{E_{d}E_{0}}{E_{0}^{2} - (h\nu)^{2}},$$
(7)

where E_d is the energy of dispersion which measures the average force of the optical transitions, E_0 is the energy of the simple oscillator. Using equations (4) and (7), we obtained the Wemple-DiDomenico model of refractive index's dispersion, which can be written as the representation $(n^2 - 1)^{-1}$ according to $(hv)^2$ on Fig. 9. After plotting the representation $(n^2 - 1)^{-1}$ according to $(hv)^2$ the linear part of the curves was extrapolated by the line y = a - bx using the linear extrapolations of each curve deduce the values from the energy of the

oscillator E_0 and dispersion E_d . The parameters a and b are related to the values of $E_0 = \left(\sqrt{\frac{a}{b}}\right)$ and $E_d = \frac{1}{\sqrt{a \cdot b}}$.

The static dielectric constant permittivity (dielectric constant at zero frequency, ε_{∞}) is defined using equation (8).

$$\varepsilon_{\infty} = \varepsilon_1(0) = \lim_{\nu \to 0} \left[n^2(h\nu) \right] = 1 + \frac{E_d}{E_0}.$$
(8)

Table 2 indicates the values for E_0 , E_d , n_0 , and ε_{∞} calculate for all chalcogenide systems. The values from dispersion energy increase while the other calculated parameters decrease when the germanium content increases.

The dissipation factor, $\tan \delta$, is described in equation (9) which is the power loss rate of the mechanical mode in a dissipative system [35].

$$\tan \delta = \frac{\varepsilon_2}{\varepsilon_1}.$$
 (9)

Figure 10 plots the dissipation factor tan δ against $h\nu$ the for Ge_xSb_{20-x}Te₈₀ (GST) alloys with composition x = 15, 17, and 19 at %. As shown in Fig. 9, tan δ values display the same variation as the dielectric constants depicted in Fig. 8. The dissipation pick factor indicates interactions between the electrons; the interactions are also the origin of the peaks in the dielectric constant spectra.



Fig. 9 - Wemple-DiDomenico plot.

We add here that the structural, optical spectroscopy, optical conductivity and dielectric properties of $BaTi_{0.5}(Fe_{0.33}W_{0.17})O_3$ and the electronic processes in noncrystalline materials are studied in [36, 37].



Fig. 10 – Spectral dependence of $(\tan \delta)$ for Ge_xSb_{20-x} Te₈₀ with composition x = 15, 17, and 19 at %.



Fig. 11 – Spectral dependence of the real and imaginary parts of optical conductivity for $Ge_xSb_{20-x}Te_{80}$ (GST) alloys with composition x = 15, 17, and 19 at %: a) $Ge_{15}Sb_5Te_{80}$; b) $Ge_{17}Sb_3Te_{80}$; c) $Ge_{19}Sb_1Te_{80}$.

Having the values of ε_1 and ε_2 we obtained the complex electric modulus (Fig. 12) and the complex impedance (Fig. 13), by follows

$$Me^* = \frac{1}{\tilde{\varepsilon}} = Me_1 + iMe_2 = \frac{\varepsilon_1}{(\varepsilon_1^2 + \varepsilon_2^2)} + i\frac{\varepsilon_2}{(\varepsilon_1^2 + \varepsilon_2^2)}, \qquad (10)$$

$$Z^{*} = \frac{1}{i\omega C_{0}\tilde{\varepsilon}} = \frac{Me^{*}}{i\omega C_{0}} = Z_{1} + iZ_{2}, \quad i^{2} = -1,$$
(11)

where Me₁ and Me₂ are the real and imaginary parts of the electric modulus, respectively, Z₁ and Z₂ are the real and imaginary parts of the complex impedance, respectively, and $C_0 = \left(\frac{A}{d}\right) \varepsilon_0$ is the vacuum capacitance of the cell, *A* and *d* is the area and the thickness of chalcogenide layers, ε_0 is the free space dielectric constant. The complex electric modulus is a decisive parameter to acquire information about the relaxation mechanism.



Fig. 12 – Variation of the complex electric modulus, Me^{*} for $Ge_xSb_{20-x}Te_{80}$ (GST) alloys with composition x = 15, 17, and 19 at %: a) Ge₁₅Sb₅Te₈₀; b) Ge₁₇Sb₃Te₈₀; c) Ge₁₉Sb₁Te₈₀.



Fig. 13 – Variation of the complex impedance, Z^* for Ge_xSb_{20-x}Te₈₀ (GST) alloys with composition x = 15, 17, and 19 at %: a) Ge₁₅Sb₅Te₈₀; b) Ge₁₇Sb₃Te₈₀; c) Ge₁₉Sb₁Te₈₀.

Figure 13 shows the variation of the real and imaginary parts of the impedance as a function of energy photon. The AFM-SEM technology from ZEISS is used to study the surface's morphology. They are revealed in Fig. 14 compact, continuous, and smooth films. The surfaces are without cracks and digs, without exfoliation area with good adhesion, and have a uniform distribution of granules originating from the film structure, influenced by the deposition conditions (material evaporation phase). The root means square (RMS) roughness of the sample surfaces measured by AFM was 1.86 nm ($Ge_{15}Sb_5Te_{80}$), 2.04 nm ($Ge_{17}Sb_3Te_{80}$;), and 1.75 nm ($Ge_{19}Sb_1Te_{80}$).



Fig. 14 – AFM and SEM images for $Ge_xSb_{20-x}Te_{80}$ (GST) alloys with composition x = 15, 17, and 19 at %: a) $Ge_{15}Sb_5Te_{80}$; b) $Ge_{17}Sb_3Te_{80}$; c) $Ge_{19}Sb_1Te_{80}$.

GST have applications in several branches of industry such as solar cells, electronic and optoelectronic devices at a low price.

This opens the door to control the structural and the optical properties of GST materials and leads to critical industrial applications such as solar cells, electronic and optoelectronic devices at a low price

4. CONCLUSIONS

The chalcogenide layers $Ge_xSb_{20-x}Te_{80}$ (GST) have been synthesized from elements with 5N purity (Ge, Sb, Te) using the conventional melt-quenching method. Structural and optical properties of GST layers are studied. The optical properties have been investigated by the Spectro-ellipsometry method. The refractive index and extinction coefficient decrease when germanium / antimony content increases/ decreases. Hight refractive index is obtained, and GST layers are commonly integrated like materials of contrasting refractivity to form transparent components with anti-reflective properties.

The optical absorption coefficient α , has shown a high value around the fundamental absorption edge and reflects the excellent quality of our samples demonstrate also by AFM and SEM imagines. The optical bandgap energy was estimated to be ~1.38 eV, which agrees with those reported by other authors. The optical dispersion parameters calculated by the single-oscillator model of Wemple-DiDomenico, show that the values from dispersion energy (E_d) increase while the others calculated parameters, the single-oscillator energy (E_0), the zero-frequency refractive index, and the dielectric constant at zero frequency, (ϵ_{∞})decrease when the germanium content increases. The dissipation factor tan(δ), the optical conductivity (σ), the complex electric modulus (M^*), and the complex impedance (Z^*) of these GST layers were calculated. The AFM-SEM imagines revealed compact, continuous, and smooth films with good adhesion and uniform granules originating from the film structure, influenced by the deposition conditions (material evaporation phase).

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